

Electromagnetism of Continuous Media — Optics

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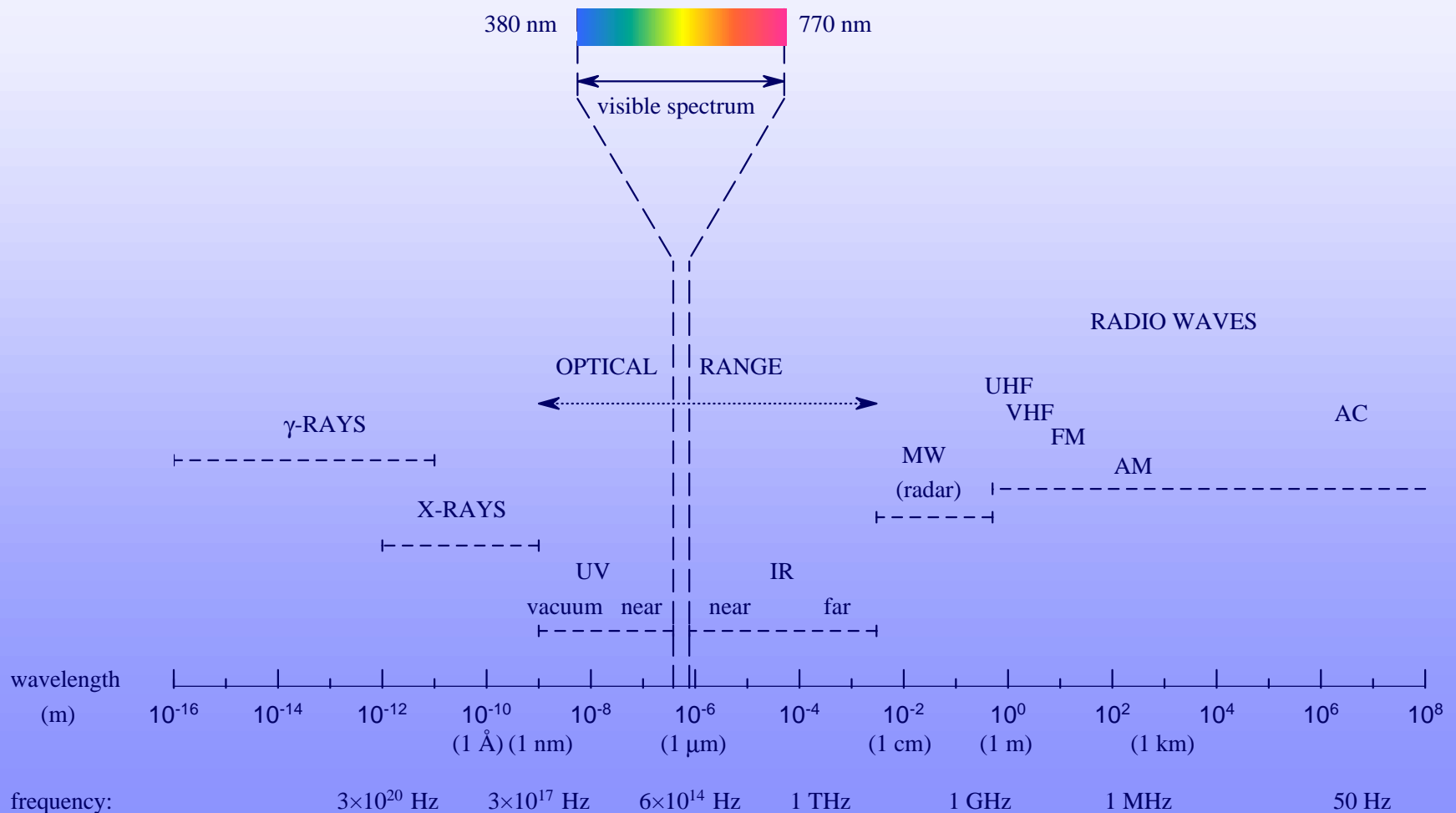
Course for 3rd year of
University studies
(Université Paris-Nord)

Series of 14 lectures in Department
of Dielectrics, IOP

Topics of Interest

- Layered systems and optical coatings: in detail
- Birefringence and polarization optics: in detail
- Dispersion and pulse propagation: yes
- Electro- and magneto-optics: yes
- Nonlinear optics: basics
- Waveguides and fibers: basics
- Inhomogeneous media: basics
- Coherence: marginal
- Diffraction, gaussian beams: no

Electromagnetic Spectrum



Notation conventions

$$\times \equiv \wedge$$

$$\mathit{curl} \mathbf{E} \equiv \mathit{rot} \mathbf{E} \equiv \nabla \times \mathbf{E} \equiv \nabla \wedge \mathbf{E}$$

$$\mathit{div} \mathbf{E} \equiv \nabla \cdot \mathbf{E}$$

Maxwell Equations

$$\nabla \wedge \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$\nabla \wedge \mathbf{H} - \frac{\partial \mathbf{D}}{\partial t} = \mathbf{j}$$

$$\nabla \cdot \mathbf{D} = \rho$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\oint \mathbf{E} \cdot d\mathbf{l} = -\frac{\partial}{\partial t} \int \mathbf{B} \cdot d\mathbf{S}$$

$$\oint \mathbf{H} \cdot d\mathbf{l} = \frac{\partial}{\partial t} \int \mathbf{D} \cdot d\mathbf{S} + J$$

$$\oint \mathbf{D} \cdot d\mathbf{S} = Q$$

$$\oint \mathbf{B} \cdot d\mathbf{S} = 0$$

$$\mathbf{E}_{2t} - \mathbf{E}_{1t} = 0$$

$$\mathbf{H}_{2t} - \mathbf{H}_{1t} = \mathbf{j}_s$$

$$D_{2n} - D_{1n} = \sigma$$

$$B_{2n} - B_{1n} = 0$$

$$\mathbf{D} = \epsilon \mathbf{E} = \epsilon_0 \mathbf{E} + \mathbf{P}$$

$$\mathbf{B} = \mu \mathbf{H} = \mu_0 \mathbf{H} + \mathbf{M}$$

Wave equation

$$\nabla^2 \mathbf{E} - \frac{N^2}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0$$

$$\nabla^2 \mathbf{H} - \frac{N^2}{c^2} \frac{\partial^2 \mathbf{H}}{\partial t^2} = 0$$

$$\mathbf{E} = \mathbf{E}_0 e^{i(\omega t - \mathbf{k} \cdot \mathbf{r})}$$

$$\mathbf{H} = \mathbf{H}_0 e^{i(\omega t - \mathbf{k} \cdot \mathbf{r})}$$

Maxwell Equations

$$\nabla \wedge \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

Faraday (induction)

$$\mathbf{E}_{2t} - \mathbf{E}_{1t} = 0$$

$$\nabla \wedge \mathbf{H} - \frac{\partial \mathbf{D}}{\partial t} = \mathbf{j}$$

Ampère (induction)

$$\mathbf{H}_{2t} - \mathbf{H}_{1t} = \mathbf{j}_s$$

Coulomb (electric charges)

$$D_{2n} - D_{1n} = \sigma$$

$$\nabla \cdot \mathbf{D} = \rho$$

No magnetic charges

$$B_{2n} - B_{1n} = 0$$

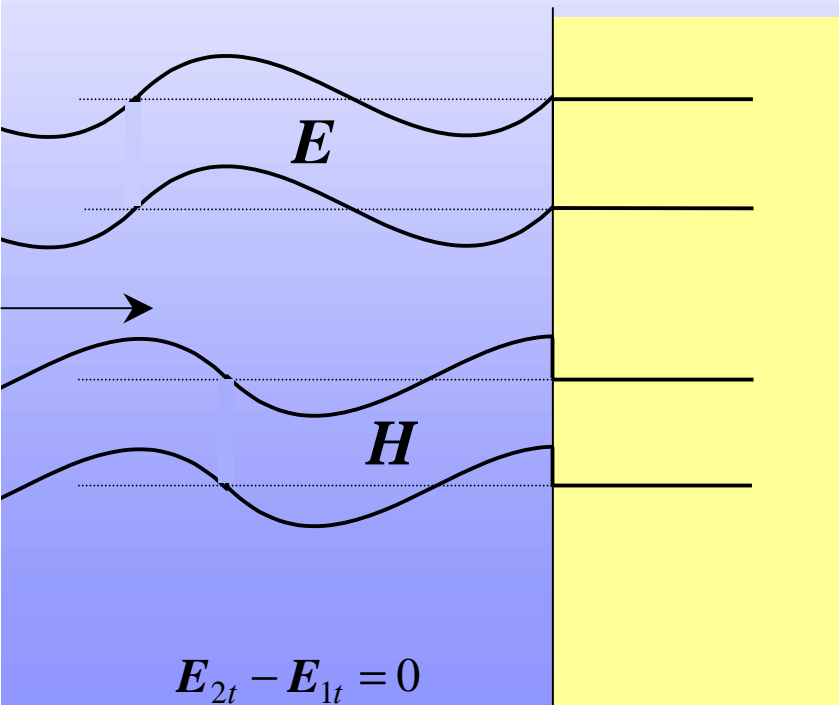
$$\nabla \cdot \mathbf{B} = 0$$

$$\mathbf{D} = \epsilon \mathbf{E} = \epsilon_0 \mathbf{E} + \mathbf{P}$$

$$\mathbf{B} = \mu \mathbf{H} = \mu_0 \mathbf{H} + \mathbf{M}$$

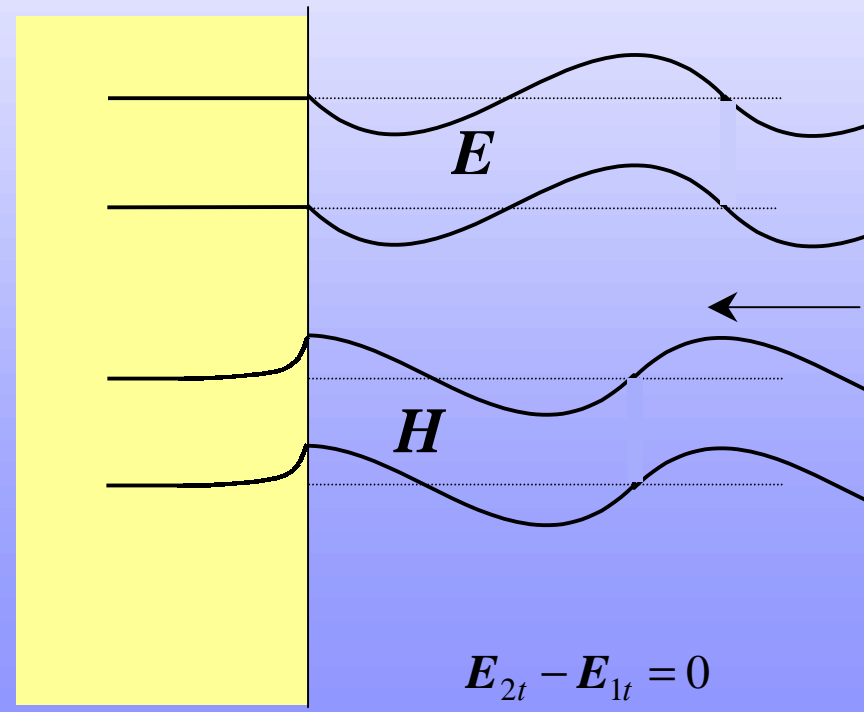
Discontinuity at the interface

- Ideal conductor ($\sigma = \infty$)



$$E_{2t} - E_{1t} = 0$$
$$H_{2t} - H_{1t} = j_s$$

- Real conductor ($\sigma < \infty$)

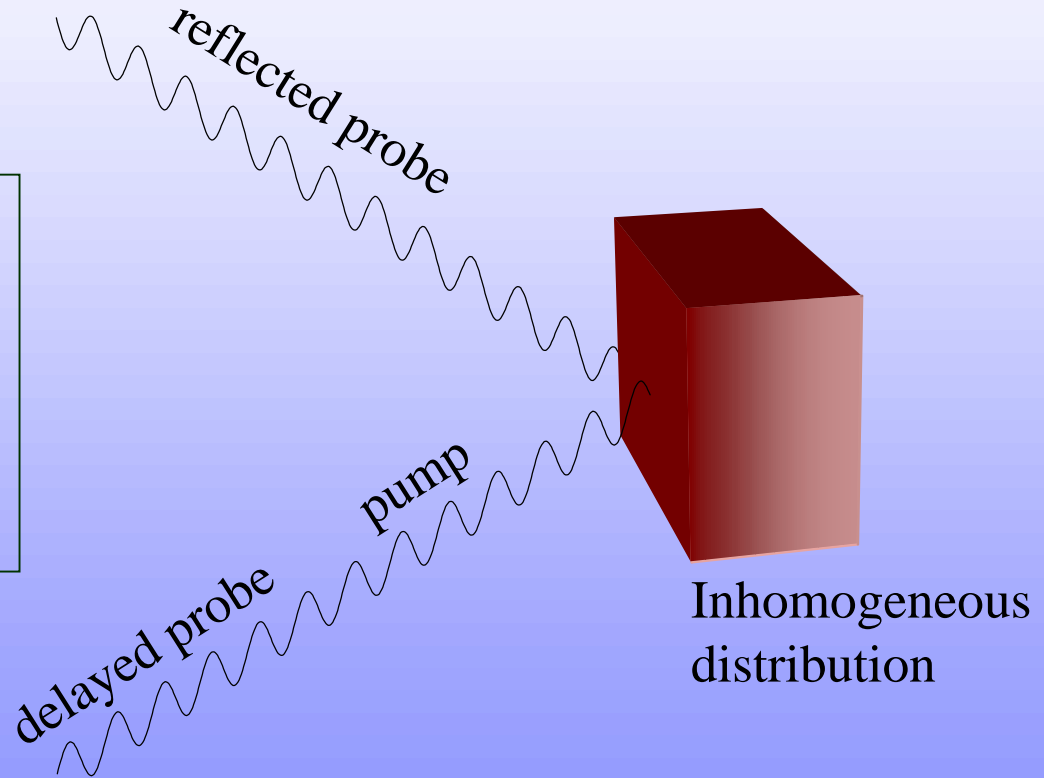


inhomogeneity

$$E_{2t} - E_{1t} = 0$$
$$H_{2t} - H_{1t} = 0$$

Time-resolved reflectivity

- Pump pulse:
high density of free carriers
(non-equilibrium state)
- Probe pulse:
monitoring of the
permittivity change



Properties of the permittivity

- dispersive [$\epsilon = \epsilon(\omega)$] — not dispersive (ϵ does not depend on ω)
- absorbing (ϵ is complex) — not absorbing or transparent (ϵ is real)
- inhomogeneous [$\epsilon = \epsilon(\mathbf{r})$] — homogeneous (ϵ does not depend on \mathbf{r})
- anisotropic (ϵ is a second rank tensor) — isotropic (ϵ is a scalar)
- nonlinear [$\epsilon = \epsilon(\mathbf{E}, \mathbf{H})$] — linear (ϵ does not depend on the fields)

$$\mathbf{D}(\omega) = \epsilon(\omega) \mathbf{E}(\omega) = \epsilon_0 \mathbf{E}(\omega) + \mathbf{P}(\omega)$$

$$\mathbf{P}(\omega) = \epsilon_0 \chi(\omega) \mathbf{E}(\omega)$$

$$\epsilon(\omega) = \epsilon_0 (1 + \chi(\omega))$$

Spectral decomposition

Real notation:

- Monochromatic wave $a(t) = |A| \cos(\omega t + \alpha) = \operatorname{Re}\{Ae^{i\omega t}\}$ with $A = |A|e^{i\alpha}$

- Polychromatic wave $a(t) = \sum_k |A_k| \cos(\omega_k t + \alpha_k)$

$$a(t) = \int_0^{\infty} |A(\omega)| \cos(\omega t + \alpha(\omega)) d\omega$$

Complex notation:

- Monochromatic wave $a(t) = Ae^{i\omega t}$ or $a(t) = \frac{1}{2}(Ae^{i\omega t} + A^*e^{-i\omega t})$

- Polychromatic wave $a(t) = \int_0^{\infty} A(\omega)e^{i\omega t} d\omega$ or $a(t) = \frac{1}{2} \int_{-\infty}^{\infty} A(\omega)e^{i\omega t} d\omega$

Calculation of the field products

- If only mean values are needed (energy density and flow):

$$\langle a(t)b(t) \rangle = \frac{1}{T} \int_0^T |A| \cos(\omega t + \alpha) |B| \cos(\omega t + \beta) dt = \frac{1}{2} |AB| \cos(\alpha - \beta)$$

in complex notation: $\langle a(t)b(t) \rangle = \frac{1}{2} \operatorname{Re}\{AB^*\}$

- If instantaneous values are needed (nonlinear optics):

$$a(t) = \frac{1}{2} (Ae^{i\omega t} + A^* e^{-i\omega t}) \qquad b(t) = \frac{1}{2} (Be^{i\omega t} + B^* e^{-i\omega t})$$

Maxwell equations for the spectral components

$$\nabla \wedge \mathbf{E} = -i\omega\mathbf{B}$$

$$\nabla \wedge \mathbf{H} = \mathbf{j} + i\omega\mathbf{D}$$

$$\nabla \cdot \mathbf{D} = \rho$$

$$\nabla \cdot \mathbf{B} = 0$$

Introduction of the generalized permittivity:

Total current (free + polarization charges):

$$\mathbf{j}_{TOT} = \sigma\mathbf{E} + i\omega\mathbf{D} = i\omega \left(\epsilon - \frac{i\sigma}{\omega} \right) \mathbf{E} = i\omega\hat{\epsilon}\mathbf{E}$$

We define a new \mathbf{D} which involves the free charges:

$$\hat{\mathbf{D}} = \hat{\epsilon}\mathbf{E}$$

Third Maxwell equation becomes:

$$\nabla \cdot \hat{\mathbf{D}} = \nabla \cdot (\epsilon - i\sigma/\omega)\mathbf{E} = \rho + \frac{1}{i\omega} \nabla \cdot \mathbf{j} = 0$$

$$\nabla \wedge \mathbf{E} = -i\omega\mathbf{B}$$

$$\nabla \wedge \mathbf{H} = i\omega\mathbf{D}$$

$$\nabla \cdot \mathbf{D} = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

Conservation of energy

All conservation laws have the general form:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{j} = P$$

Work of the field per unit of time exerted on the charges (Lorentz force):

$$q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B}) \cdot \mathbf{v} = q\mathbf{v} \cdot \mathbf{E} \rightarrow \mathbf{j} \cdot \mathbf{E}$$

From Maxwell equations:

$$\left. \begin{array}{l} \nabla \wedge \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \\ \nabla \wedge \mathbf{H} = \mathbf{j} + \frac{\partial \mathbf{D}}{\partial t} \end{array} \right\} \begin{array}{l} \cdot \mathbf{H} \\ \cdot \mathbf{E} \end{array} \left. \vphantom{\begin{array}{l} \nabla \wedge \mathbf{E} \\ \nabla \wedge \mathbf{H} \end{array}} \right\} \mathbf{E} \cdot \frac{\partial \mathbf{D}}{\partial t} + \mathbf{H} \cdot \frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{E} \wedge \mathbf{H}) = -\mathbf{j} \cdot \mathbf{E}$$

$$\frac{\partial U}{\partial t} + \nabla \cdot \mathbf{S} = -\mathbf{j} \cdot \mathbf{E}$$

Conservation of energy: continued

$$\mathbf{E} \cdot \frac{\partial \mathbf{D}}{\partial t} + \mathbf{H} \cdot \frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{E} \wedge \mathbf{H}) = -\mathbf{j} \cdot \mathbf{E} \longleftrightarrow \frac{\partial U}{\partial t} + \nabla \cdot \mathbf{S} = -\mathbf{j} \cdot \mathbf{E}$$

Linear non-dispersive (and non-absorbing) media:

$$U = \frac{1}{2} (\mathbf{E} \cdot \mathbf{D} + \mathbf{H} \cdot \mathbf{B})$$

Absorbing medium: for a monochromatic wave the energy change corresponds to the heat:

$$Q = \frac{\omega}{2} \text{Im}(\epsilon^*) \mathbf{E} \mathbf{E}^* = \omega \text{Im}(\epsilon^*) \langle E^2 \rangle$$

Medium with small dispersion: quasi-monochromatic wave:

$$\langle U \rangle = \frac{1}{4} \left(\frac{d(\omega \epsilon)}{d\omega} \mathbf{E} \cdot \mathbf{E}^* + \mu_0 \mathbf{H} \cdot \mathbf{H}^* \right) = \frac{1}{2} \left(\frac{d(\omega \epsilon)}{d\omega} \langle E^2 \rangle + \mu_0 \langle H^2 \rangle \right)$$