

Magnetic properties of $(\text{Nd}_{0.33}\text{Eu}_{0.33}\text{Gd}_{0.33})\text{Ba}_2\text{Cu}_3\text{O}_y$ doped by additional Gd-211 particles

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Abstract: Magnetic hysteresis loops were measured on a series of $(\text{Nd}_{0.33}\text{Eu}_{0.33}\text{Gd}_{0.33})\text{-Ba}_2\text{Cu}_3\text{O}_y$ (NEG) samples melt-grown in a reduced oxygen atmosphere with additions of different concentrations of $\text{Gd}_2\text{BaCuO}_5$ (Gd-211) particles. Measurements were conducted by means of a SQUID magnetometer in fields up to 7 T, $5 \leq T \leq 92$ K. The hysteresis loops exhibit a strong peak effect in a wide range of temperatures. Scaling properties were studied and interpreted in terms of a phenomenological model for RE-123 materials. It is shown that a dispersion of fine Gd-211 particles is very effective in increasing the critical current density.

Key words: Current densities, pinning forces, scaling, $(\text{Nd}_{0.33}\text{Eu}_{0.33}\text{Gd}_{0.33})\text{Ba}_2\text{Cu}_3\text{O}_y$ superconductors

INTRODUCTION

The embedding of small normal-conducting particles into melt-textured superconductors is way to increase the flux pinning strength. Depending on the size of the particles, this pinning is mainly effective at low temperatures/low fields. The newly developed NEG superconductors were found to exhibit a strongly developed peak (fishtail) effect [1]. To increase the critical current density, J_c , even further, we added (Nd,Eu,Gd)-211 particles with the same ratio as in the NEG-123 matrix. By TEM analysis, it was found that we indeed obtain submicron-sized 211 particles, but the particles consist mainly of Gd in the rare earth site [2]. Therefore, we attempt the addition of various amounts of Gd-211 to the starting composition.

EXPERIMENTAL

Powders of Nd_2O_3 , Eu_2O_3 , Gd_2O_3 , BaCO_3 and CuO were weighed to have a nominal composition of $(\text{Nd}_{0.33}\text{Eu}_{0.33}\text{Gd}_{0.33})\text{Ba}_2\text{Cu}_3\text{O}_y$ (NEG). These powders were ground thoroughly and calcinated at 880 °C for 24 h with intermediate grinding, which was repeated three times, and were then pressed into pellets and sintered at 1020 °C for 48 h. NEG bulk samples with a volume fraction of 10 (A), 20 (B), 30 (C) and 40 mol% (D) of Gd-211 secondary phase were prepared using a mixture of sintered NEG and commercial Gd-211 powders. Since the Pt addition is known to be effective in refining the size of Y-211 in

the YBCO system, we added 0.5 mol% Pt to the starting powders. More details of the sample preparation can be found in Refs. [3, 4] Finally, the samples are cut into a cubic shape with dimensions $1 \times 1.5 \times 0.3 \text{ mm}^3$. $T_{c,\text{onset}}$ of all the samples is around 93.5 K. Magnetization loops are obtained using a Quantum Design MPMS-7 SQUID magnetometer equipped with a 7 T superconducting magnet; $H_a \parallel c$. The homogeneity of all samples was checked by magneto-optic imaging [5].

RESULTS AND DISCUSSION

The field dependence of the critical current density, $J_c(B)$ for samples A, B, C, and D, is presented in Fig. 1, measured at $T = 77 \text{ K}$. All samples show a pronounced fishtail (peak) effect. The position of the fishtail peak is always found at $\approx 2.5 \text{ T}$. Only sample D shows

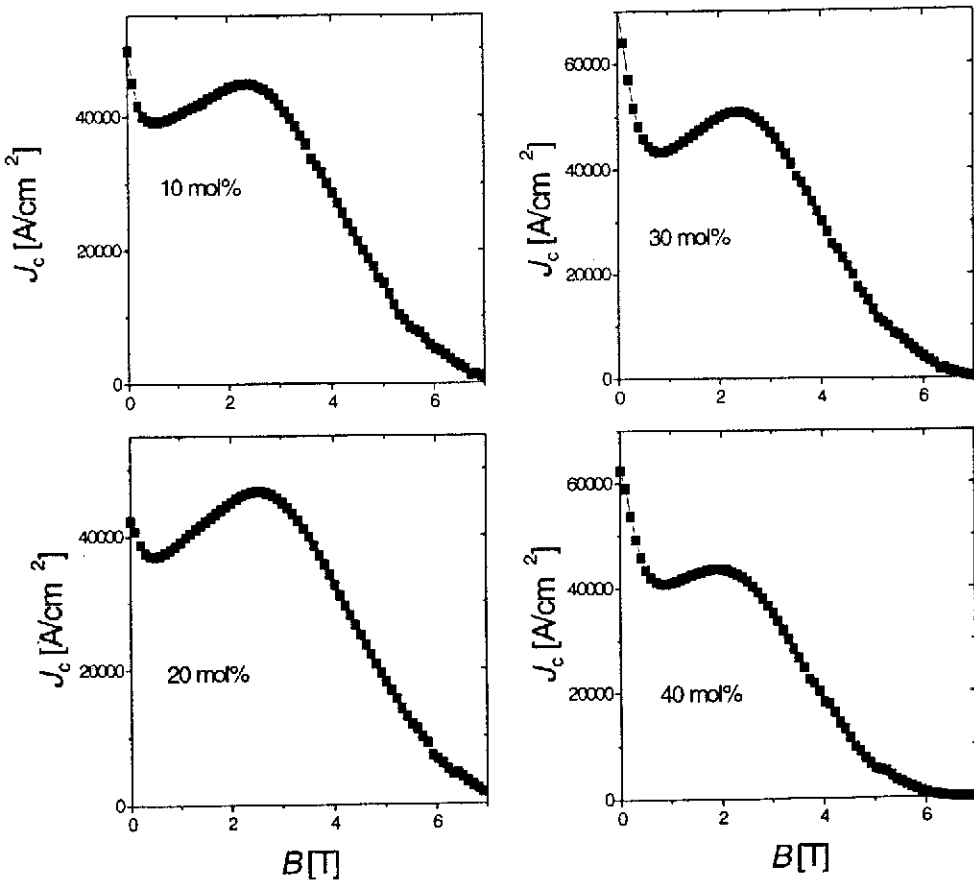


Fig. 1. $J_c(B)$ of all samples A – D with different Gd-211 content, measured at $T = 77 \text{ K}$.

a somewhat reduced peak field. The current density, $J_c(0)$, is increasing on increasing Gd-211 content. This demonstrates that the submicron-sized Gd-211 particles [2] are effective to increase the pinning at low fields, but an addition of a too high amount of Gd-211

particles as in the case of sample D clearly affects the properties of the superconducting NEG matrix. Note that J_c is largest for sample C. Therefore, an addition of 30 mol% Gd-211 particles seems to be the optimum, as shown in Ref. [2].

To analyze the pinning properties of superconductors, the scaling of the current densities, J_c , or volume pinning forces, $F_p = J_c \times B$, plays an important role [6, 7, 8]. In Ref. [9], an empirical formula was developed

$$f(b_f) = b_f^{m+1} \exp\left[\frac{m+1}{n}(1 - b_f^n)\right], \quad (1)$$

with $b_f = B_a/B_{fp}$ and $f = F_p/F_{p,max}$ are the applied field and the pinning force density normalized with respect to the coordinates of the fishtail maximum (B_{fp} , $F_{p,max}$). Jirsa and Púst [9] showed that eq. (1) is capable to describe the scaling of various 123 superconductors found in literature with $m = 1$ and leaving n as the only free parameter. For a fixed position of the fishtail peak n can be interpreted as a measure of the width of the normalized fishtail curve. Alternatively, n represents a measure of the mutual position of B_p and B_{irr} , where B_{irr} denotes the irreversibility field. The higher n , the lower the ratio B_{irr}/B_p , i.e. a high n indicates that the fishtail peak lies close to irreversibility field. It is now important to relate the parameters m and n with microstructural features. In fig. 2, we present the scaling of samples A and D, where we expect to observe differences as revealed from the current density measurements (fig. 1).

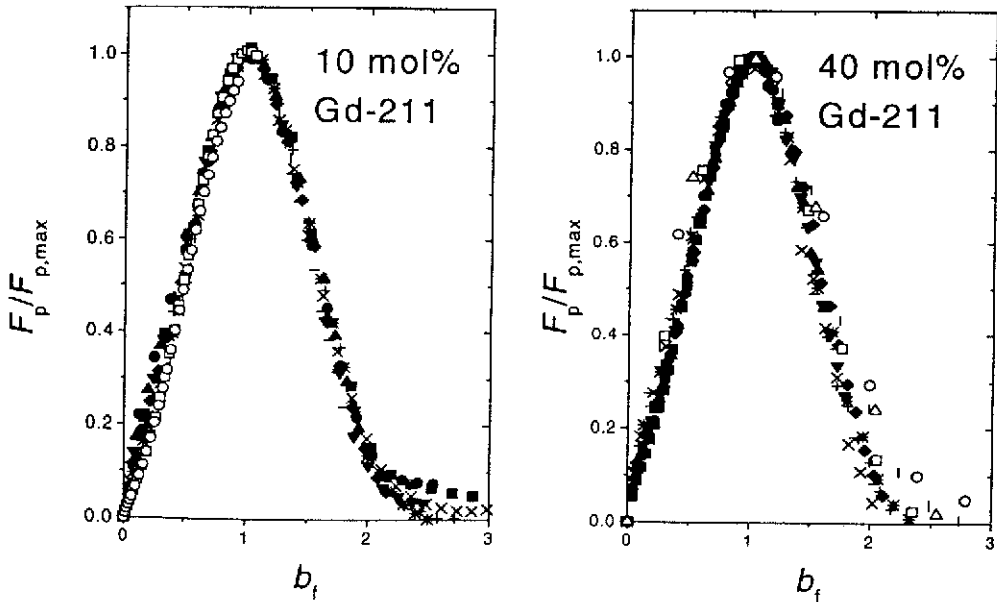


Fig. 2. F_p -scaling of sample A and D at temperatures between 60 K and 90 K. The scaling is found to be excellent; in sample D some deviations from the ideal behaviour occur.

The scaling was performed using data between 60 K and 90 K, where all necessary parameters are directly accessible to experiment. The scaling of both samples is excellent; the

parameter n is determined to 2.34 (sample A) and 2.13 (sample D). However, the fitting of sample D reveals some important differences: The data at low b_f increase faster than the best fit to eq. (1). This indicates that at low fields, there is a strong contribution to the pinning by the normal conducting particles. The data taken at temperatures above 84 K also show a deviation at the low field side. As discussed in Ref. [1], the main source of pinning in the NEG compounds is of the δT_c -type. These pinning sites may become normal at elevated temperatures. The shape change of the $j_c(B)$ curves and the deviation in the scaling of sample D may therefore indicate that the dominant pinning mechanism changes from δT_c to δl -type (i.e. pinning at normal-conducting particles).

CONCLUSIONS

The addition of Gd-211 particles is effective to increase the pinning strength at low fields: the optimum content is found to be 30 mol%. The scaling of the volume pinning forces is found to be excellent, with n ranging between 2.13 and 2.34.

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